

Fluctuations of grains inside a discharging two-dimensional silo

Angel Garcimartín,* Iker Zuriguel, Alvaro Janda, and Diego Maza

Departamento de Física, Facultad de Ciencias, Universidad de Navarra, E-31080 Pamplona, Spain

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We present experimental data corresponding to a two-dimensional dense granular flow, namely, the gravity-driven discharge of grains from a small opening in a silo. We study the local velocity field at the scale of single grains at different places with the help of particle-tracking techniques. From these data, the velocity profiles can be obtained and the validity of some long-standing approaches can be assessed. Moreover, the fluctuations of the velocities are taken into consideration to characterize the features of the advective motion (due to the gravity force) and the diffusive motion, which shows nontrivial behavior.

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I. INTRODUCTION

The process of silo discharge can be considered a canonical example of the complexity displayed by granular flows [1]. Typically, when an orifice is opened at the bottom of a container filled with grains, the particles will exit through it and the outflow gives rise to a velocity field inside the silo. For big orifices the flow will quickly homogenize, while for small orifices it will develop intermittencies [2]. A clear-cut definition allowing to qualify what *big* or *small* really means is unavailable, but it is normally accepted that an orifice diameter of about five times the typical size of the beads can be taken as a boundary between both situations. Recent systematic works suggest that even larger orifices are needed to avoid jamming events and intermittencies [3–5]. Regardless of these two regimes, the prediction of the flow rate at the outlet remains an open problem, at least in the sense that a derivation of a quantitative expression from fundamental principles has not been worked out yet. Simple dimensional analysis [6] provides a relationship between the orifice size and the mass flow, but a theoretical foundation is lacking to support it. One of the reasons for the absence of this theoretical framework is that the flow near the outlet is accelerated due to the breakdown of the balance between gravity and dissipation. Nevertheless, far from the exit orifice (“far” meaning a few orifice diameters above it), the stationary velocity field is well defined if the outlet size is big enough. This suggests that the velocity profile could be described as a function of the stress field. However, an equation relating both variables has proven elusive. Perhaps the most important reason for this is the absence of a suitable spatial scale over which to average variables such as the density or the velocity.

Two utterly different approaches have been employed to predict the velocity profile inside a silo. The first one, profusely used in soil mechanics, starts from the consideration of the granular material as a continuous deformable media, and leads to concepts like yield stress and plastic potentials [7]. Despite its success in generally describing many situations, some concerns remain about its applicability whenever discontinuities in the stress and velocity fields appear, which are not unusual. The second approach assumes from the very beginning the discrete nature of the material. The collective movement is

taken as the sum of individual particles. It is easier to consider the diffusion of “voids” instead of particles; these voids enter in the silo through the outlet orifice and move upward [8]. This notion is certainly suggestive and, once developed, predicts qualitatively well the velocity profile [9]. Unfortunately, it fails to match quantitatively the experimental measurements of the velocity.

Nedermann and Tüzün [10] proposed the simple argument that the horizontal velocity component u should be proportional to the gradient of the vertical velocity component v along the horizontal direction, that is,

$$u = -B \frac{\partial v}{\partial x}, \quad (1)$$

where B is a constant with length dimensions. Hence, assuming mass conservation and the incompressibility of the flowing material, it is easy to obtain a parabolic Partial Differential Equation (PDE) for the spatial dependence of v

$$\frac{\partial v}{\partial y} = -B \frac{\partial^2 v}{\partial x^2}. \quad (2)$$

Note that this expression is valid only after a transient flow regime, when a steady state is reached. As Eq. (2) is isomorph to the diffusion equation, the constant B is usually referred to as a “diffusion length,” although such a statement is not rigorously true, given that B was introduced just as an *ad hoc* constant in Eq. (1). From these premises, and assuming that the size of the orifice is negligible compared to the silo dimensions, the velocity profile in the vertical direction is

$$v = -\frac{Q}{\sqrt{4\pi B y}} \exp\left(-\frac{x^2}{4B y}\right), \quad (3)$$

where Q is the volumetric flow rate. Notably, this expression has been successfully used to fit some experimental observations with an acceptable accuracy. The main shortcoming of this formula when confronted with the measurements is the wide discrepancy among the values of B reported by different authors [11]. In some cases, a dependence of B with the vertical coordinate has also been noted [12].

Efforts have been recently made to find some common ground among all those lines of reasoning. As the concept of diffusive motion of individual particles or the holes left by them is broadly accepted in this context, some authors have

*angel@fisica.unav.es

generalized this notion by proposing the collective motion of groups of particles [13]. These authors put forward the existence of “spots,” which are zones where the vacant space left by the absence of a particle is shared by a group of grains, causing a local decrease of the packing fraction. Such spots would diffuse upward. In addition, they introduce a new length scale in the problem which allows to solve the discrepancies between the values of B measured in the experiments and those predicted by the old diffusive approach. If diffusion is mediated by spots, the displacements of the particles are not necessarily of about their own size; moreover, the scale of the spot displacement can be *quantitatively* related to the constant B introduced by kinetic arguments.

No matter how valuable this idea is, the microscopic foundations of the velocity profile and the mass flow rate remain unclear. Diffusive models are certainly inspiring but do not describe correctly important features of the flow, such as the fluctuations of the velocity field or the mean square displacement of the particles at short scales. These magnitudes have been considered in a series of recent papers reporting anomalous Probability Density Function (PDF) for the velocity distribution around its mean values [14–16]. Also, the mean square displacements reported suggest that at short times the particle rearrangements are not ballistic and therefore the global behavior of the material cannot be regarded as a classical hydrodynamic flow. Although the accuracy of some of the cited references represent a qualitative leap compared to earlier works, the experimental conditions used to obtain them prevent a direct comparison with theoretical models. For instance, it is sometimes argued that polydisperse grains are used to avoid crystallization effects; nevertheless, the disorder introduced by this practice cannot be easily taken into account by any existing microscopic model. Besides, many of these experiments are performed in a quasi-two-dimensional configuration that could display, in fact, three-dimensional effects, and the influence of this dimension on the flow properties is unclear. Finally, let us mention that recent approaches suggested the possibility of understanding the particle dynamics by means of elastoplastic hybrid ideas [17] or a high density kinetic theory [18].

In this work we attempt to avoid the mentioned troubles by performing our experiment in a single layer of particles. Although a strict two-dimensional (2D) configuration is experimentally unaccessible, we designed a silo where the dynamics is driven only by the rearrangements of the particles within a single layer, and the effect of the front and rear walls just reduces to an extra dissipative term. We also use monodisperse particles, which indeed tend to crystallize, but as we will show in the following, both the microscopic and the macroscopic behavior of the granular flow are compatible with the above-mentioned descriptions. We restrict ourselves to an orifice that is big enough to guarantee that the jamming probability is negligible.

The paper is organized as follows. We will first describe the experimental setup. Then we will show and comment the measurements corresponding to the velocity profiles. A section follows where we will focus on the velocities of individual particles. Then we will describe the fluctuations, and finally we will gather some conclusions.

II. EXPERIMENTAL PROCEDURES AND METHODS

As the experiment aim is the study of the movement of individual particles, we have built a two-dimensional silo made of two glass sheets. Its dimensions are 800-mm high and 200-mm wide. Between the two panes a thin frame of stainless steel (1.1-mm thick) is sandwiched. This frame fixes the gap between the two glass sheets, and at the same time it contains the beads at the sides (see Fig. 1). At the bottom, there is an orifice, its size being settled at 15.9 mm. Although this length is somewhat arbitrary, it is a compromise. On the one hand, the opening is well above the limit for jamming [3–5], and therefore we have never observed an obstruction of the exit. On the other hand, the experiment is less demanding if the flow velocity is not too high, and as the outflow depends on the size of the opening, it cannot be arbitrarily large.

The granular sample consists of monodisperse stainless steel beads with a diameter $d = 1.00 \pm 0.01$ mm. The grains are poured along the whole width of the silo from a hopper at the top. The granular deposit within the silo consists of a monolayer of particles because the gap between the glass sheets is slightly larger than the beads. Then a small overlap between the spheres in the images is possible. In our device, however, the overlaps detected are smaller than about 2% of the particle diameter. The overlap is likely to be even smaller if one considers that flowing beads are not necessarily in contact. Finite size effects in the lateral dimensions (corresponding to the width of the recipient) can be neglected as the silo dimensions (width and height) are much bigger than the size of the beads and the width of the opening at the bottom. The

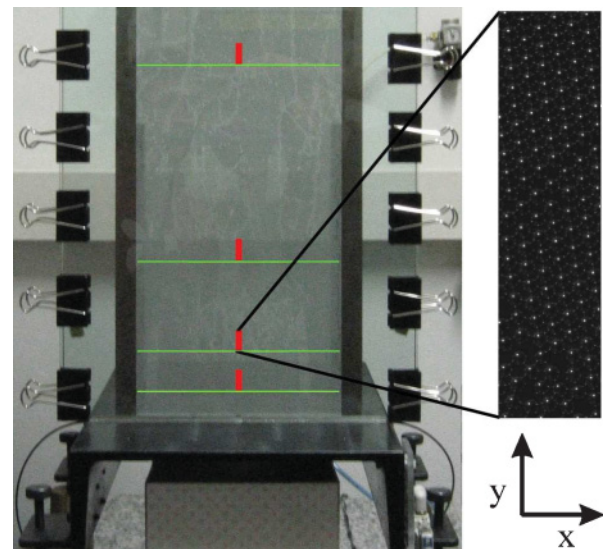


FIG. 1. (Color online) Sketch of the silo. The photograph covers approximately the lower half of the silo. The two vertical strips at the sides are steel sheets that contain the beads and set the separation between the front and rear glass panes. The rectangular red patches show the four zones where images were recorded, which are described in the text. A typical image from one of the films is shown at the side. Horizontal green lines show the places where the measurement of the velocity profile was performed (see text). The choice of coordinates is also shown. The origin of the coordinates coincides with the center of the exit orifice.

same setup was used in previous works [5,19]; those previous results were taken as benchmarks for some of the data provided in this article.

Recordings were obtained with a high speed camera (Photron model FASTCAM 1024 PCI). With suitable optics and lighting, a small zone of the silo can be videotaped with high resolution. We will now describe these zones where measurements have been performed, displayed in Fig. 1. Four heights were chosen, viz. $y = 30, 65, 155,$ and 355 mm (y is measured from the bottom of the silo). A region of interest was then defined as a rectangle of 256×1024 pixels (width by height), the bottom side of the rectangle placed at the above-mentioned heights. The spatial resolution is of 50 pixels per millimeter, so the rectangle corresponds to about 5×20 mm. These four zones are centered in the vertical axis of the silo and are marked with rectangular patches in Fig. 1. Under convenient illumination, each bead reflects a bright spot, as seen in the same figure, where a single frame is shown as an example. Films of about 3 seconds were recorded of the four zones at a speed ranging from 2000 to 5400 frames per second (the recording speed is changed so as to optimize the displacement of the beads from frame to frame). The length of the films was limited by the memory available in the recording system. We noticed that the steady state regime for the flow can be reached as quickly as in a few seconds, and we waited for about 10 seconds after opening the exit orifice before recording the films.

To rule out the possibility that the observation zone somehow influenced the results, we also recorded a second set of films changing a little the region of interest. In this second set, the window was a square of 512×512 pixels, with the bottom side at the same height than the previous rectangles. The recording speed was also changed by a small amount. As no differences were perceived in the results, in this paper we will just show data corresponding to one set of recordings.

Once the films were recorded and stored, we processed them with a particle tracking software that was written in our laboratory to deal with a similar situation [20]. The procedure involves splicing the film into frames, detecting the spots reflected by the beads, and linking up the positions of the beads in adjacent frames. As the beads move less than one radius from one frame to the next, this can be done unambiguously. The software can pinpoint the position of the beads with subpixel resolution; the accuracy of the procedure has been checked with stationary beads, yielding a figure of ± 0.05 pixels. In this way, we have been able to obtain about 1000 bead tracks for each film (order of magnitude).

Another series of recordings was performed to obtain the vertical velocity v along a horizontal line. With this aim, we recorded a set of films with a much lower magnification (about 5 pixels per millimeter), of a horizontal band (1024×64 pixels), spanning almost the whole 200 mm of the layer and centered at the vertical axis. The colored horizontal lines shown in Fig. 1 indicate those regions. The speed of the recording was much smaller in these cases (500 or 1000 frames per second usually). The bands where measurements are made are about 12.8-mm high, so there is a small gradient of v inside the measuring zone. We can evaluate this after Eq. (3); at $x = 0$, the relative variation of the vertical velocity

is $\frac{\Delta v}{v} = \frac{\Delta y}{2y}$. If we put $\Delta y = 12.8/2$, we obtain that in the worst case ($y = 30$), $\frac{\Delta v}{v} \approx 10\%$ at the middle of the measuring zone.

The procedure carried out in this case was the following. A series of short films (typically 24 recordings, each one lasting about 0.1 second) were registered. Then for each particle, the mean vertical velocity was calculated. This amounts to fitting a straight line to the tracks and calculating the slope. All the data were aggregated into a single file for each zone. About 10^5 individual particles (order of magnitude) were detected and processed at each height. For completeness, the mean velocity field for the whole silo, albeit with much smaller resolution, obtained with a Particle Image Velocimetry (PIV) technique, can be seen in Ref. [5].

III. VELOCITY PROFILE ALONG THE HORIZONTAL DIMENSION

As explained in the former section we have carried out the measurements of v , spanning all the width of the silo, at four heights. When preparing an experimental realization, it can be observed with the naked eye that particles tend to group in large clusters with preferred orientation order; the frontiers between them are clearly visible. It is not surprising, then, that the particles move orderly, displaying groups of data points with similar velocities (Fig. 2). To obtain suitable statistical averages, close to 10^5 particle velocities were obtained at each height. Moreover, these measurements are the accumulation of more than 20 different runs, taken at different times. It is only after performing such a large amount of measurements that averages fit neatly to Eq. (3). In Fig. 2 we have plotted the average velocities along with the Gaussian fit. As remarked in the preceding section, the velocity is slightly different inside the measuring zone. To confirm that this effect is small, we have performed some measurements in the upper and lower layers inside the measuring zone. For instance, we measured the velocities in the top and bottom 3 mm of the zone at $y = 30$ (recall that this is the worst case), and obtained a peak velocity at $x = 0$ of 153 and 133 mm/s, respectively, which is to be compared with the mean value of 142 mm/s [Fig. 2(d)]. Therefore, it can be concluded that although this effect is present, it does not account for the dispersion of the velocities as displayed by the clouds of gray points in Fig. 2.

Although Eq. (3) is strictly valid only for a *point-like* orifice at the bottom, we used it to fit the experimental observations leaving two free parameters B and Q . We performed the fit with the method of least squares; the values obtained are shown in Table I. The values of Q are in reasonable agreement with the measured flow rate ($Q = 4600$ beads per second). This quantity was independently measured by weighing the material poured from the silo during a fixed amount of time. An alternative procedure could have been to fix Q and leave B as the only free parameter to be fitted by the least-squares method (the results are similar). In any case, it is noteworthy that the values of B are close to 2, which was the initial guess of Nedderman and other authors [7]. We do not observe a neat dependence of B with height except for the upper region, but this could be due to other factors. The velocities are much smaller, and the region where beads move downward reach the lateral walls, so the boundaries could influence the velocity field.

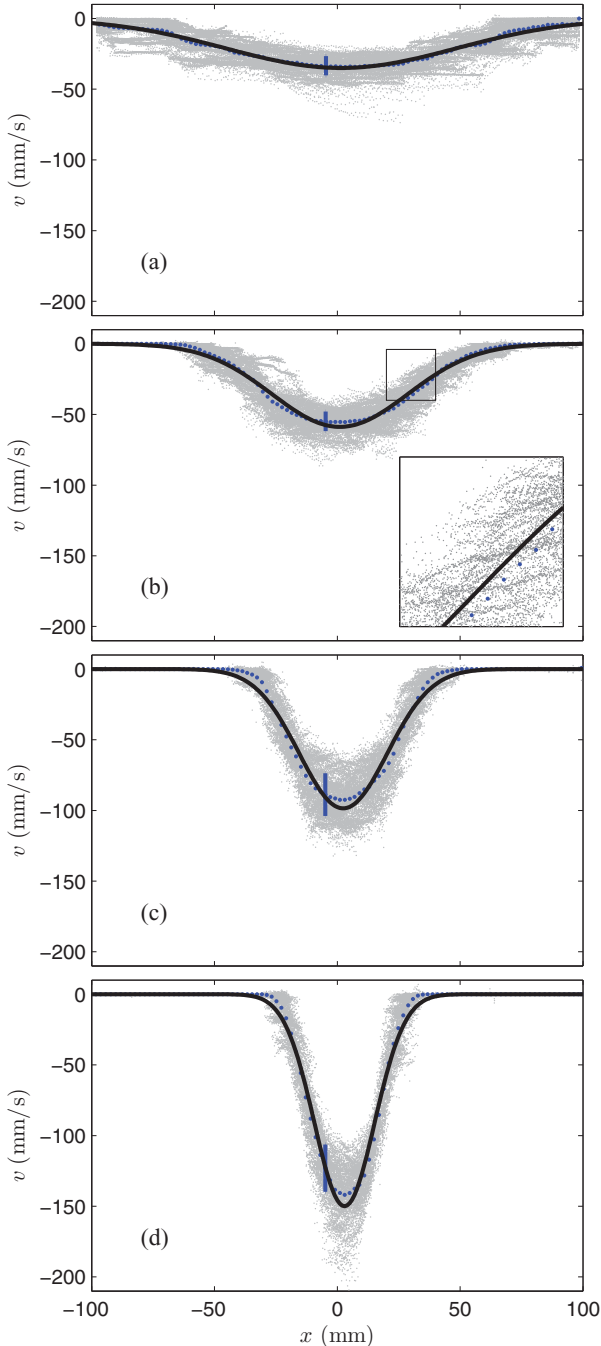


FIG. 2. (Color online) Vertical component of the velocity at different heights. The measurements were taken at the horizontal lines depicted on Fig. 1: (a) $y = 355$, (b) $y = 155$, (c) $y = 65$, (d) $y = 30$ mm. The cloud of small gray points corresponds to the velocities of individual particles. Blue dots correspond to the local mean (obtained every 2 mm, approx.) A bar is provided displaying the standard deviation of the measurements around the mean (to avoid cluttering the figure just one bar is shown at $x \approx -5$ mm). Black solid lines correspond to the fit of Eq. (3). Note that all the axes are at the same scale. The inset in (b) is a zoom of the zone indicated with a frame.

As remarked earlier, the agreement between the experimental data and Eq. (3) is good for the average velocities. However, the deviations of the velocities of individual particles from the

TABLE I. Values of the parameters B and Q [Eq. (3)] as given by the fit shown in Fig. 2.

y (mm)	B (mm)	Q (beads per second)
355	3.07	4100
155	2.46	4100
65	2.46	4400
30	2.50	4600

average are quite big, as revealed by the standard deviation. Moreover, the collective motion of grains can be discerned if one looks closely at the data points, as shown in the inset of Fig. 2(b), where the aligned clouds of points show these events. This means that the motion of particles is not just noisy, but also that the grains move in *clusters*. It is therefore pertinent to take a closer look at the movement of individual grains.

IV. VELOCITIES OF INDIVIDUAL PARTICLES

We will now focus our attention on the path of individual particles resolved in time, and we will describe it through the instantaneous velocity of each particle. As explained in Sec. II, the frame rate of the recording was tuned to obtain a point at a sampling frequency such that the displacement of the particles from frame to frame is close to the spatial resolution of the system. As the velocity implies a derivation, if one performs a straightforward calculation the noise is too large. To avoid it, we have calculated the velocity averaged at a length scale corresponding to a bead radius. The method is the following. The track of one particle is taken, and from each point of the track the location in the trace at a distance equal to the bead radius is searched. The time elapsed from the current point to that position is calculated by interpolation and from this time lapse the velocity was computed as the radius divided by the elapsed time. This figure yields the speed (velocity modulus), but as the direction of the bead movement can also be obtained, we next extract the vertical component v . In the following we will deal exclusively with the absolute value of this component. (Remark, however, that the horizontal component of the velocity is much smaller than the vertical one, so the results we will present do not change noticeably if the speed is considered instead of v). In fact, the method explained amounts to low-pass filtering the vertical velocity at a length scale of a bead radius. This procedure was carried out with more than 1000 beads for each one of the four rectangular patches depicted in Fig. 1. The results are shown in Fig. 3.

The average vertical velocity of all the particles in the observation region at a given instant, which we will call $\langle v \rangle$, can be calculated from the velocities of individual particles. It is also represented in the left column graphics of Fig. 3 with the solid black lines. Then the mean vertical velocity V , averaged over time, is obtained; it is represented in the same plots with a dashed horizontal line. The values obtained in the four observation zones, from top to bottom, are $V = 37.0, 56.6, 85.1,$ and 144.4 mm/s. This compares well to the values provided by the fit of Eq. (3), which are $V = 35.0, 58.8, 98.6,$ and 150 mm/s, respectively, at $x = 0$. As seen in Fig. 2, the measured velocities at the center are a little bit smaller than those given by the fit: the maxima are (in absolute value)

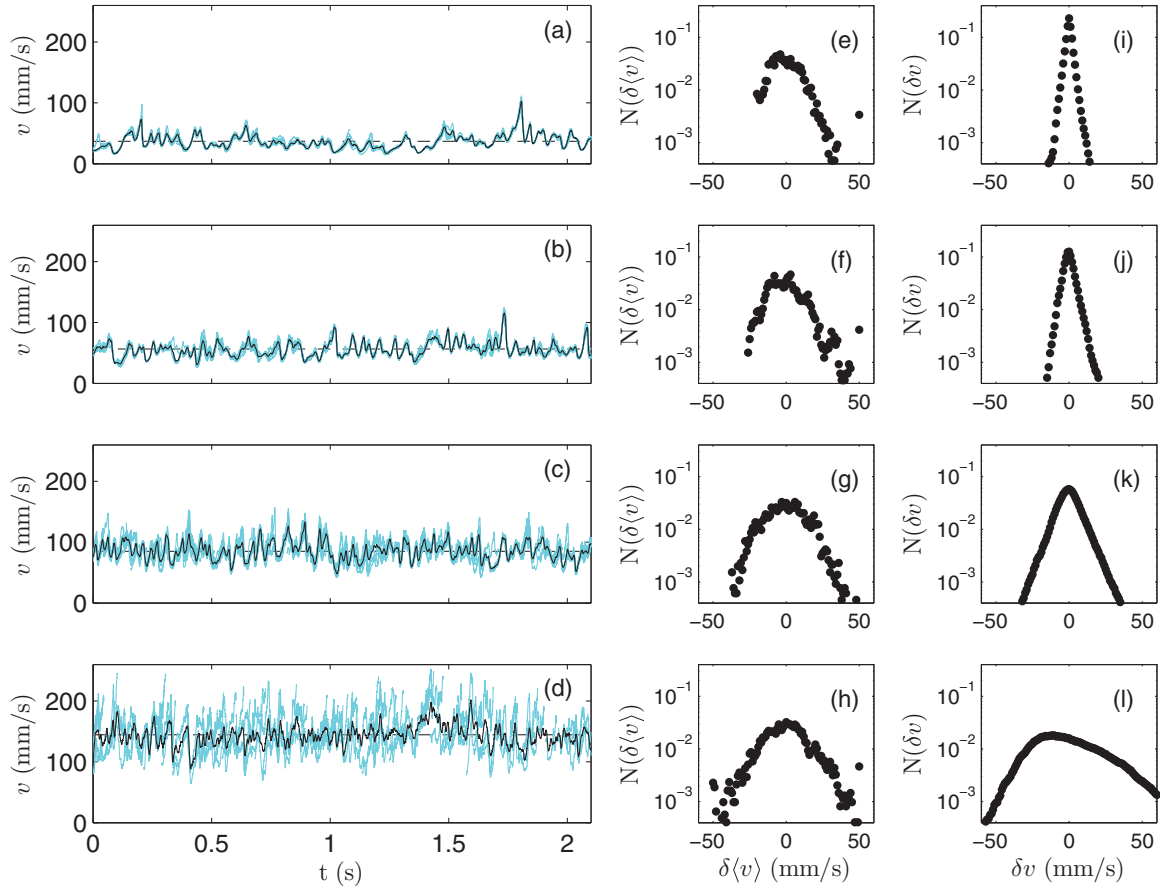


FIG. 3. (Color online) On the left column (a)–(d), the vertical component of the velocity of individual particles v , obtained as explained in the text, is shown with cyan (gray) lines. The four plots correspond to the rectangular regions depicted in Fig. 1: (a) $y = 355$, (b) $y = 155$, (c) $y = 65$, (d) $y = 30$ mm. Only 5% of the traces are represented to avoid cluttering the graphics. The solid black lines are the instantaneous averages $\langle v \rangle$, taking in this case all the measured particles, not just the ones represented here. The dashed horizontal line corresponds to the mean vertical velocity V averaged over all the time for all the beads. Note that all the axes are at the same scale to allow easy comparison. The plots on the center column (e)–(h) display the normalized histograms of the fluctuations of $\langle v \rangle$ (the average velocities represented in the left column with a solid black line) around the mean V . Each plot corresponds to the same zone as the one to its left. Note the logarithmic scale on the vertical axis. All the axes are at the same scale to allow easy comparison. The right column (i)–(l) corresponds to the normalized histograms of the fluctuations of the individual velocities v around the average $\langle v \rangle$. Each plot corresponds to the same zone as the one to its left. Note the logarithmic scale on the vertical axis. All the axes are at the same scale, which is also the same as in the center column, to allow easy comparison.

$V = 34.3, 55.4, 92.5$, and 142 mm/s, respectively. In any case, the small time recorded and shown in Fig. 3 can be considered representative of the typical velocity values.

A direct observation of the leftmost column of Fig. 3 indicates that the fluctuations of the average speed $\langle v \rangle$ increase as the orifice is approached. But not only that, the deviations of individual particles from the average also seem to be enhanced. It can be appreciated with the naked eye that at the top of the silo particles move in concert, as a group, closely assembled around the mean; at the bottom, the dispersion around the average seems more important. This inspired us to separate the fluctuations into two: the variation of the average itself and the deviations of individual particles from the average. This is not a standard procedure (for instance, when studying turbulent fluids) because in the usual situations there is a well defined time averaged velocity $\langle v \rangle$ that does not fluctuate wildly over time; in our case this is not so. In the following section we will take a closer look into these fluctuations.

V. FLUCTUATIONS

The fluctuations of the velocities are thus separated into two categories: the fluctuations of the average velocity $\langle v \rangle$ around V and the fluctuations of individual velocities around the average $\langle v \rangle$. The former can be regarded as the fluctuations of a certain advective velocity field (due to the gravity), while the latter can be assigned to the random movement of the particles around the advective field.

The fluctuations of $\langle v \rangle$ are shown on the center column of Fig. 3. Although they are not exactly Gaussian, one can discern a shape resembling a parabola (in semilogarithmic scale). This probably means that the average velocity is just noisy and that the fluctuations arise from a stochastic process. In fact, the autocorrelation function of $\langle v \rangle$ (Fig. 4) shows that it is uncorrelated for times longer than the delay that it takes for a bead to fall its own diameter. A similar behavior was found in numerical simulations [21]. An important remark concerning

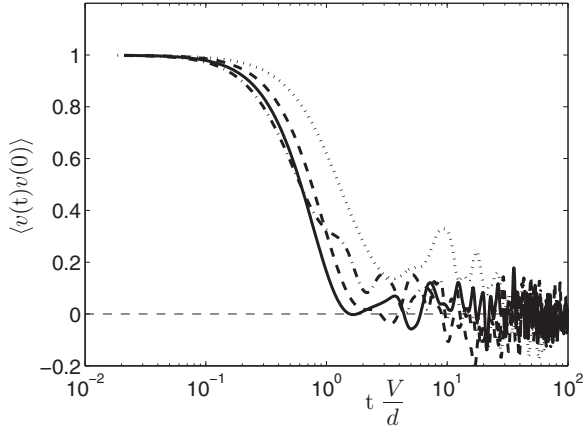


FIG. 4. Autocorrelation of the average velocity $\langle v \rangle$. The horizontal axis has been rescaled by the average time it takes for a bead to fall its own diameter. The four lines correspond to the four heights where measurements were performed: dot-dashed line, $y = 355$; solid line, $y = 155$; dashed line, $y = 65$, and dotted line, $y = 30$ mm.

this should be noted: the fact that the autocorrelation of $\langle v \rangle$ goes to 1 for short times was something to be expected from the method we used to calculate the velocities, which erases the fluctuations at times smaller than $t \frac{V}{d}$. The relevant result here is that the autocorrelation decays to zero for times larger than this. In Fig. 4 the rescaled time lag is shown to stress the point that at about $t = \frac{V}{d}$ the average velocity $\langle v \rangle$ quickly becomes uncorrelated. In such a representation, the autocorrelation decays later at the bottom of the silo. But let us note that, in fact, the *unrescaled* decay times are larger where the velocity is smaller (i.e, at the top of the silo, the autocorrelation of $\langle v \rangle$ decays later than at the bottom).

On the other hand, the fluctuations of v around $\langle v \rangle$ are clearly not Gaussian (right column of Fig. 3). They are pointed and sharp, meaning that the deviations from the average are *smaller* than if they were the result of a random process. This is a hint of a collective, correlated motion of the particles. This feature is much more remarkable far from the exit orifice, at the top of the silo, where the mean velocity V is small. At the bottom, the distribution is wider and less sharp. This again agrees with previous numerical simulations [22].

The comparison of both sets of fluctuations (the deviations of $\langle v \rangle$ around V and the deviations of the individual velocities v around $\langle v \rangle$) reveals that their amplitudes are of the same order of magnitude (Fig. 3). Another common feature to both sets of fluctuations, which is in good agreement with previous observations, is that the distributions are slightly slanted to the right. That is, there are slightly more large deviations toward the region of higher velocities (meaning falling more rapidly) than toward the region of smaller velocities. As mentioned above, the main difference between the PDFs is that while the deviations of $\langle v \rangle$ around V are approximately Gaussian, the deviations of v around $\langle v \rangle$ display a pointed shape revealing collective motion. This suggests an underlying process of anomalous diffusion.

To investigate this, we have computed the variance of the displacements of the beads. In this calculation the “mean path” must be subtracted from the track of each particle.

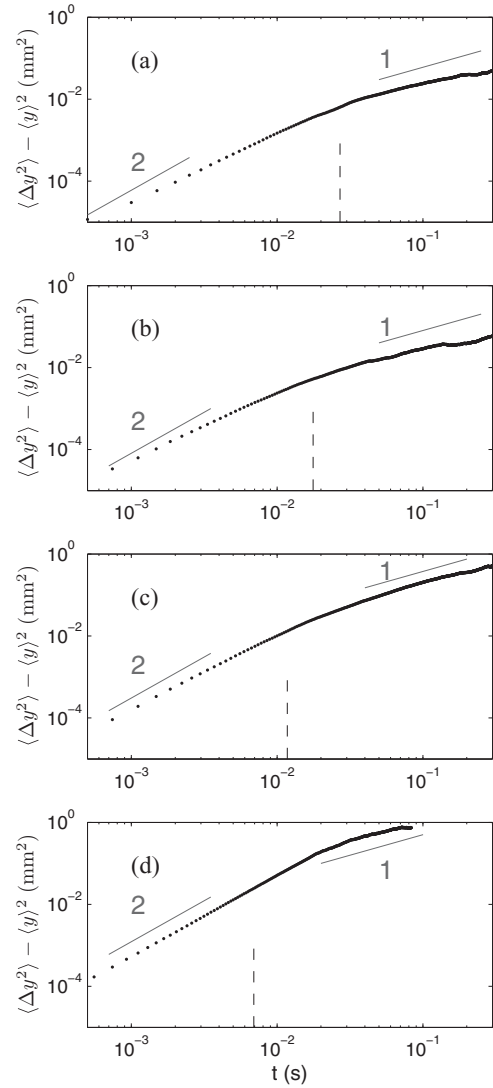


FIG. 5. The variance of the displacements of the beads, calculated as explained in the text, at the four heights (a) $y = 355$, (b) $y = 155$, (c) $y = 65$, and (d) $y = 30$ mm. Note that all the axes are at the same scale. The solid lines are a guide for the eye: in each plot the line at the left has a slope of 2 and the one at the right a slope of 1, as labeled. The dashed vertical lines indicate the instant at which the particle has covered, on average, a distance equal to its own diameter.

Calculating directly the mean position of a group of beads can be tricky because particles entering or exiting the observation window must be explicitly taken into account. Instead, we have integrated the average velocity $\langle v \rangle$ obtained previously to calculate the mean displacement. This quantity at each instant is then subtracted from the particle tracks to compute the variance of the vertical displacements $\sigma_y^2 = \langle (y - \langle y \rangle)^2 \rangle$. The same method can also be performed for the x component, but as the results do not add new information we will only present the variance for the y component (Fig. 5). In particular, an explicit comparison is drawn with a diffusive regime ($\sigma_y^2 \sim t$) and with a ballistic regime ($\sigma_y^2 \sim t^2$). It has been pointed out, with the support of experimental data [23], that fluctuations of dense grains should generically display a ballistic regime for small times. Then, as time increases, a regime might

develop showing the cage effect [24], in which basically the fluctuations remain at the same level with time (i.e., $\sim t^0$). Finally, a diffusive regime is expected for long enough times. Times can be gauged in terms of the distance traveled by the particle. Small times are those smaller than the interval it takes for a bead to move a distance equal to its own diameter; the cage effect, if it exists, and the transition to a diffusive regime, should be more conspicuous from there on. In Fig. 5 we have marked the moment at which the particle has traveled a distance equal to its diameter (on average).

The picture that emerges is similar to that of Brownian motion resolved for very small times [25], except that an intermediate stage occurs where the particle is “trapped in a cage” by its neighbors. Our data are not able to display explicitly this scenario because the time that one can continuously follow a particle is limited by the observation window, or else by the gradients inside the observation window, if one tries to consider bigger and bigger windows. We can nevertheless identify a region for small t where the fluctuations scale approximately as t^2 . This scaling evolves in time and for displacements larger than a diameter the fluctuations appear to scale with time subdiffusively [i.e., the slope is smaller than 1 (Fig. 5)]. Remarkably, this evolution is particularly evident for the region near the orifice [Fig. 5(d)] where the ballistic regime extends beyond the limit of one particle diameter. In this region the particles seem to be in a less dense state, where displacements are mainly determined by the gravity field that accelerates the particles and causes the collisions, which is the mechanism governing the diffusive motion.

VI. DISCUSSION

We provided experimental data about the microscopic movement of particles in a 2D dense granular flow. The spatial and time resolutions allowed us to explore the collective motion. It has been shown that the averaged vertical velocity profile can be fitted to a Gaussian function which is a solution of the diffusive-like equation. Although the instantaneous velocity of the particles display correlated motions, long time averages show the typical velocity profile reported in different experimental situations [14,15,18]. It has been sometimes

argued that systems consisting of monodisperse particles cannot be used to study granular flows due to their tendency to crystallize. Nevertheless, these results show that even in this case the diffusive approach can be used to fit the average velocity profile. Importantly, the election of monodisperse beads in a 2D configuration gives the opportunity to test the fundamental assumption of the theoretical models related to the existence of a well-defined temporal scale of the particle displacement at a microscopic level. As we show in Sec. V the particle motion can be split into a global or *advective* movement and another one “on top” of it where the diffusive approach can be considered. The former displays a commonplace Gaussian PDF of the velocity fluctuations while the latter shows a strong non-Gaussian PDF, conspicuously indicating collective motion. The Gaussian-like fluctuations seem to indicate the randomness of the global deformation of the granular layer. On the other hand, the evolution of the autocorrelation function of the averaged velocity $\langle v \rangle$ in the vertical direction (Fig. 4) shows that the correlation of this variable becomes negligible at the time scale of one particle displacement, regardless of the mean velocity V . Finally, the variance of the path of individual particles with respect to the mean path displays a ballistic regime for small displacements, suggesting that the global dynamics is governed by gravity at these scales. The anomalous regimes reported by other authors [14] could presumably be related to the combination of individual and collective motion. In addition, it seems that the variance at long time scales is anomalous, which could be related to the intrinsic features of the advective field.

In summary, the study of the deviations of individual velocities v from the instantaneous average $\langle v \rangle$ has revealed an inspiration to clarify the origin of the microscopic dynamics in a silo discharge process.

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